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## COMMENT

# Intertwiner realization of a simple non-standard $\boldsymbol{R}$-matrix $\dagger$ 

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Received 19 March 1993, in final form 26 July 1993


#### Abstract

Using Jimbo's method, we realize the non-standard $R$-matrix $\breve{R}^{\frac{1}{4}}(x, y)$ as an intertwiner between non-generic irreducible representations of the quantum affine algebra $U_{q}\left(\hat{s} l_{2}\right)$.


The intertwiner theory [1] developed by Jimbo is powerful in constructing solutions to the Yang-Baxter equation (YBE) with spectral parameters [2-4]. It turns out that the so-called standard $R$-matrices with spectral parameters, for example, the $R$-matrix associated with the six-vertex model, can be realized as the intertwiners between two parametrized irreducible representations of the quantum universal enveloping algebra $U_{q}(g)$ [5,6] with $g$ being an affine Lie algebra. Then the question naturally arises whether we can also put the non-standard $R$-matrices [7] into the framework of this theory to obtain a unified handling of solutions to Ybe. In this comment we try to answer this question through a simple example, in which we explain the non-standard $R$-matrix $\mathscr{R}^{\frac{1}{2}}(x, y)$ as an intertwiner between two parametrized irreducible representations of $U_{q}\left(\hat{s} l_{2}\right)$ at roots of unity.

Definition. The quantum affine algebra $U_{q}\left(\hat{s} l_{2}\right)$ is an associative algebra over the complex number field $\mathbb{C}$ generated by the elements $e_{i}, f_{i}, h_{i}(i=0,1)$ and the unit 1 subject to the following relations:

$$
\begin{array}{lrl}
{\left[h_{i}, e_{i}\right]=2 e_{i}} & {\left[h_{i}, f_{i}\right]=-2 f_{i}} \\
{\left[h_{i}, e_{j}\right]=-2 e_{j}} & {\left[h_{i}, f_{j}\right]=2 f_{j}} & (i \neq j) \\
{\left[h_{0}, h_{1}\right]=0} & {\left[e_{i}, f_{j}\right]=\delta_{i j}\left[h_{i}\right] \equiv \delta_{i j}\left(q^{h_{i}}-q^{-h_{i}}\right) /\left(q-q^{-1}\right)}  \tag{1}\\
e_{i}^{3} e_{j}-[3] e_{i}^{2} e_{j} e_{i}+[3] e_{i} e_{j} e_{i}^{2}-e_{j} e_{i}^{3}=0 \\
f_{i}^{3} f_{j}-[3] f_{i}^{2} f_{j} f_{i}+[3] f_{i} f_{j} f_{i}^{2}-f_{j} f_{i}^{3}=0 & (i \neq j) .
\end{array}
$$

It is well known that $U_{q}\left(\hat{s} l_{2}\right)$ can be endowed with the coproduct $\Delta$

$$
\begin{aligned}
& \Delta\left(e_{i}\right)=q^{h_{i}} \otimes e_{i}+e_{i} \otimes 1 \\
& \Delta\left(f_{i}\right)=f_{i} \otimes q^{-h_{i}}+1 \otimes f_{i} \\
& \Delta\left(h_{i}\right)=h_{i} \otimes 1+1 \otimes h_{i}
\end{aligned}
$$

[^0]and the antipode $S$
$$
S\left(h_{2}\right)=-h_{i} \quad S\left(e_{i}\right)=-q^{-h_{1}} e_{1} \quad S\left(f_{2}\right)=-f_{i} q^{h_{1}}
$$
to become a Hopf algebra, and according to Jimbo's argument for $x \in \mathbb{C}-\{0\}$ there exists a homomorphism of algebras $\varphi_{x} \cdot U_{q}\left(\hat{s} l_{2}\right) \rightarrow U_{q}\left(s l_{2}\right)$ given by
\[

$$
\begin{array}{lll}
\varphi_{x}\left(e_{0}\right)=x f & \varphi_{x}\left(f_{0}\right)=x^{-1} e & \varphi_{x}\left(h_{0}\right)=-h \\
\varphi_{x}\left(e_{1}\right)=e & \varphi_{x}\left(f_{1}\right)=f & \varphi_{x}\left(h_{1}\right)=h
\end{array}
$$
\]

where $e, f$ and $h$ are the generators of $U_{q}\left(s l_{2}\right)$, which satisfy

$$
[h, e]=2 e \quad[h, f]=-2 f \quad[e, f]=[h]
$$

Then from a representation $(\pi, V)$ of $U_{q}\left(s l_{2}\right)$, where $V$ is the representation space and $\pi$ a homomorphic mapping from $U_{q}\left(s l_{2}\right)$ to $\operatorname{End}(V)$, one can form the composition $\left(\pi \cdot \varphi_{x} \otimes \pi \cdot \varphi_{y}\right) \cdot \Delta$.

$$
U_{q}\left(\hat{s} l_{2}\right) \xrightarrow{\Delta} U_{q}\left(\hat{s} l_{2}\right) \otimes U_{q}\left(\hat{s} l_{2}\right) \xrightarrow{\pi \cdot \varphi_{1} \otimes \pi \cdot \varphi_{r}}(\text { End } V \otimes \text { End } V)=\operatorname{End}(V \otimes V)
$$

which gives rise to a representation of $U_{q}\left(\hat{s} l_{2}\right)$ depending on $x, y \in \mathbb{C}-\{0\}$.
Jimbo's remarkable result [1] says that the standard $R$-matrices associated with $U_{q}\left(s l_{2}\right)$ are intertwiners (module isomorphisms) between certain representations $\left(\left(\pi \cdot \varphi_{x} \otimes \pi \cdot \varphi_{y}\right) \cdot \Delta, V \otimes V\right)$ and $\left(\left(\pi \cdot \varphi_{y} \otimes \pi \cdot \varphi_{x}\right) \cdot \Delta, V \otimes V\right)$ oin $U_{q}\left(\hat{s} l_{2}\right)$ with $q$ being generic.

On the other hand, it has been proved that the non-standard $R$-matrices without spectral parameters associated with $U_{q}\left(s l_{2}\right)$ can be obtained from the universal $R$-matrix by considering the representations of $U_{q}\left(s l_{2}\right)$ at roots of unity [8]. So to realize the non-standard $R$-matrices with spectral parameters as intertwiners between irreducible representations of $U_{q}\left(\hat{s} l_{2}\right)$ we naturally consider the representations of $U_{q}\left(\hat{s} l_{2}\right)$ in the case that $q$ is a root of unity.

Let us focus our attention on the simplest case that $q^{2}=-1$, in which one has a two-dimensional irreducible representation of $U_{q}\left(s l_{2}\right)$ depending on an arbitrary parameter $\lambda \in \mathbb{C}$. Suppose the representation space $V$ is spanned by the vectors $v_{0}$ and $v_{1}$, then the representation ( $\pi, V)$ can be written as (for simplicity, from now on we will use module language) [8]:

$$
\begin{array}{lrll}
h v_{0}=\lambda v_{0} & h v_{1}=(\lambda+2) v_{1} & e v_{0}=v_{1} & e v_{1}=0 \\
f v_{0}=0 & f v_{1}=-[\lambda] v_{0} & &
\end{array}
$$

and the representation $\left(\left(\pi \cdot \varphi_{x} \otimes \pi \cdot \varphi_{J}\right) \cdot \Delta, V \otimes V\right)$ of $U_{q}\left(\hat{s} l_{2}\right)$ takes the following form

$$
\begin{aligned}
& h_{0}\left(v_{i} \otimes v_{j}\right)=-(2 \lambda+2 i+2 j)\left(v_{i} \otimes v_{j}\right) \\
& h_{1}\left(v_{i} \otimes v_{j}\right)=(2 \lambda+2 i+2 j)\left(v_{i} \otimes v_{j}\right) \quad i, j=0, i \\
& e_{0}\left(v_{0} \otimes v_{0}\right)=0 \quad e_{0}\left(v_{0} \otimes v_{1}\right)=-q^{-\lambda} y[\lambda] v_{0} \otimes v_{0} \\
& e_{0}\left(v_{1} \otimes v_{0}\right)=-x[\lambda] v_{0} \otimes v_{0} \\
& e_{0}\left(v_{1} \otimes v_{1}\right)=-x[\lambda] v_{0} \otimes v_{1}+q^{-\lambda} y[\lambda] v_{1} \otimes v_{0} \\
& f_{0}\left(v_{0} \otimes v_{0}\right)=y^{-1} v_{0} \otimes v_{1}+q^{\lambda} x^{-1} v_{1} \otimes v_{0} \\
& f_{0}\left(v_{0} \otimes v_{1}\right)=-q^{\lambda} x^{-1} v_{1} \otimes v_{1}
\end{aligned}
$$

$$
\begin{aligned}
& f_{0}\left(v_{1} \otimes v_{0}\right)=y^{-1} v_{1} \otimes v_{1} \quad f_{0}\left(v_{1} \otimes v_{1}\right)=0 \\
& e_{1}\left(v_{0} \otimes v_{0}\right)=v_{1} \otimes v_{0}+q^{\lambda} v_{0} \otimes v_{1} \\
& e_{1}\left(v_{0} \otimes v_{1}\right)=v_{1} \otimes v_{1} \quad e_{1}\left(v_{1} \otimes v_{0}\right)=-q^{\lambda} v_{1} \otimes v_{1} \\
& e_{1}\left(v_{1} \otimes v_{1}\right)=0 \\
& f_{1}\left(v_{0} \otimes v_{0}\right)=0 \quad f_{1}\left(v_{0} \otimes v_{1}\right)=-[\lambda] v_{0} \otimes v_{0} \\
& f_{1}\left(v_{1} \otimes v_{0}\right)=-q^{-\lambda}[\lambda] v_{0} \otimes v_{0} \\
& f_{1}\left(v_{1} \otimes v_{1}\right)=-[\lambda] v_{1} \otimes v_{0}+[\lambda] q^{-\lambda} v_{0} \otimes v_{1} .
\end{aligned}
$$

Using these equations one can easily prove that as a $U_{q}\left(\hat{s} l_{2}\right)$ module the vector space $V \otimes V$ is generated by the vector $v_{1} \otimes v_{1}$ when $x / y \neq q^{2 \lambda}$ and $[2 \lambda] \neq 0$. In fact, from the vector $v_{1} \otimes v_{1}$ we can obtain the following vectors through the actions of $U_{q}\left(s l_{2}\right)$ :

$$
\begin{aligned}
& e_{0} v_{1} \otimes v_{1}=-x[\lambda] v_{0} \otimes v_{1}+q^{-\lambda} y[\lambda] v_{1} \otimes v_{0} \\
& f_{1} e_{0} v_{1} \otimes v_{1}=\left(x-q^{-2 \lambda} y\right)[\lambda]^{2} v_{0} \otimes v_{0} \\
& f_{0} f_{1} e_{0} v_{1} \otimes v_{1}=\left(x-q^{-2 \lambda} y\right)[\lambda]^{2}\left(y^{-1} v_{0} \otimes v_{1}+q^{\lambda} x^{-1} v_{1} \otimes v_{0}\right)
\end{aligned}
$$

and it is easily seen that when $x / y \neq q^{-2 \lambda}$ and $[2 \lambda] \neq 0$ these vectors together with the vector $v_{1} \otimes v_{1}$ span the whole space $V \otimes V$.

Proposition $\lambda$. If $x / y \neq q^{ \pm 2 \lambda}$ and $[2 \lambda] \neq 0$, the representation $\left(\left(\pi \cdot \varphi_{x} \otimes \pi \cdot \varphi_{y}\right) \cdot \Delta, V \otimes V\right)$ is irreducible.

Proof. Suppose $S$ is a non-empty $U_{q}\left(\hat{s}_{2}\right)$-invariant subspace of $V \otimes V$, then an element $w \in S$ can be written as

$$
w=\sum_{i, j=0}^{1} c_{i j} v_{i} \otimes v_{j} \quad c_{i j} \in \mathbb{C}
$$

Since there exist only the following three cases, the proposition directly follows from the fact that $V \otimes V$ is generated by $v_{1} \otimes v_{1}$.

Case 1. $c_{00} \neq 0$ or $c_{00}=0, c_{01}-q^{\lambda} c_{10} \neq 0$.

$$
\begin{aligned}
& e_{1} w=c_{00}\left(v_{1} \otimes v_{0}+q^{\lambda} v_{0} \otimes v_{1}\right)+\left(c_{01}-q^{\lambda} c_{10}\right) v_{1} \otimes v_{1} \\
& f_{0} e_{1} w=c_{00}\left(y^{-1}-x^{-1} q^{2 \lambda}\right) v_{1} \otimes v_{1} .
\end{aligned}
$$

Case 2. $c_{00}=0$ and $c_{01}=q^{\lambda} c_{10} \neq 0$.

$$
\begin{aligned}
f_{0} w & =-x^{-1} q^{\lambda} c_{01} v_{1} \otimes v_{1}+y^{-1} c_{10} v_{1} \otimes v_{1} \\
& =\left(y^{-1}-x^{-1} q^{2 \lambda}\right) c_{10} v_{1} \otimes v_{1} .
\end{aligned}
$$

Case 3. $c_{00}=0$ and $c_{01}=q^{\lambda} c_{10}=0$.

$$
w=c_{11} v_{1} \otimes v_{1}
$$

For the relation between the representations $\left(\left(\pi \cdot \varphi_{x} \otimes \pi \cdot \varphi_{y}\right) \cdot \Delta, V \otimes V\right)$ and $\left(\left(\pi \cdot \varphi_{y} \otimes \pi \cdot \varphi_{x}\right) \cdot \Delta, V \otimes V\right)$ we have

Proposition 2. If $x / y \neq q^{ \pm 2 \lambda}$ and $q^{2 \lambda} \neq 1$, there exists an intertwiner between the representations $\left(\left(\pi \cdot \varphi_{x} \otimes \pi \cdot \varphi_{y}\right) \cdot \Delta, V \otimes V\right)$ and $\left(\left(\pi \cdot \varphi_{y} \otimes \pi \cdot \varphi_{x}\right) \cdot \Delta, V \otimes V\right)$.

Proof. We introduce the notation

$$
\phi \equiv\left(\pi \cdot \varphi_{x} \otimes \pi \cdot \varphi_{y}\right) \cdot \Delta \quad \psi \equiv\left(\pi \cdot \varphi_{y} \otimes \pi \cdot \varphi_{x}\right) \cdot \Delta
$$

then what we need to prove is that there is an automorphism $\check{R}(x, y)$ of the vector space $V \otimes V$ such that the diagram

is commutative. To this end let us consider the linear mapping $\check{R}\left(x, y^{\prime}\right)$ determined by the equations

$$
\begin{aligned}
& \breve{R}(x, y) v_{0} \otimes v_{0}=v_{0} \otimes v_{0} \\
& \check{R}(x, y) v_{0} \otimes v_{1}=\left(1 /\left(y q^{\lambda}-x q^{-\lambda}\right)\right)\left(\left(q^{\lambda}-q^{-\lambda}\right) y v_{0} \otimes v_{1}+(y-x) v_{1} \otimes v_{0}\right) \\
& \check{R}(x, y) v_{1} \otimes v_{0}=\left(1 /\left(q^{\lambda}-x q^{-\lambda}\right)\right)\left((y-x) v_{0} \otimes v_{1}+x\left(q^{\lambda}-q^{-\lambda}\right) v_{1} \otimes v_{0}\right) \\
& \check{R}(x, y) v_{1} \otimes v_{1}=\left(x q^{\lambda}-y q^{-\lambda}\right) /\left(y q^{\lambda}-x q^{-\lambda}\right) v_{1} \otimes v_{1} .
\end{aligned}
$$

After some calculation one can easily see that when $x / y \neq q^{ \pm 2 \lambda}, q^{2 \lambda} \neq 1 \check{R}(x, y)$ is an automorphism of $V \otimes V$, and its commutativity with the action of $U_{q}\left(\hat{s} \hat{l}_{2}\right)$ can also be verified directly. This proves the proposition.

Written in matrix form, the intertwiner $\check{R}(x, y)$ is
$\check{R}(x, y)=\frac{1}{y t-x t^{-1}}\left[\begin{array}{llll}y t-x t^{-1} & & & \\ & \left(t-t^{-1}\right) y & (y-x) & \\ & (y-x) & \left(t-t^{-1}\right) x & \\ & & & x t-y t^{-1}\end{array}\right] \quad t=q^{\lambda}$.
This is exactly the so-called non-standard $R$-matrix with spectral parameters associated with the fundamental representation of $U_{q}\left(s l_{2}\right)$. We have successfully realized it as an intertwiner between irreducible representations of $U_{q}\left(\hat{s} l_{2}\right)$ (in the case that [2 $\left.\lambda\right] \neq 0$ ) and it seems reasonable to expect that the other non-standard $R$-matrices with spectral parameters can be handled similarly.

Finally, we should mention that the $R$-matrix obtained above can also serve as an intertwiner of certain representations of the quantum superalgebra $U_{q}(\hat{s} l(1 \mid 1))[9]$. We should also mention that it has already been pointed out in [10] that the same $R$-matrix can be understood as the intertwiner of representations of the quantum affine algebra at $q$ roots of unity. But a further explanation is not presented there. So compared with [10], this present paper includes some new results. We have treated the four-dimensional tensor representations in a mathematically rigorous way and we have derived the irreducibility condition precisely, which is not at all self-evident when $q$ is a root of unity. Besides, we have made it clear that in the non-generic case the parameter $t$ in the intertwiner is not the same as the deformation parameter $q$ in the quantum affine algebra, in contrast to the generic case.

## Acknowledgment

The author thanks Professor Min Qian for helpful discussions and thanks the referee for bringing [9] and [10] to his attention.

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[^0]:    $\dagger$ This work is supported by the National Natural Science Foundation of China.
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